

A Lyman Break Galaxy in the Epoch of Reionization from HST Grism Spectroscopy

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ABSTRACT

Slitless grism spectroscopy from space offers dramatic advantages for studying high redshift galaxies: high spatial resolution to match the compact sizes of the targets, a dark and uniform sky background, and simultaneous observation over fields ranging from five square arcminutes (*HST*) to over 1000 square arcminutes (*Euclid*). Here we present observations of a galaxy at $z = 6.57$ — the end of the reionization epoch — identified using slitless *HST* grism spectra from the PEARS survey (Probing Evolution And Reionization Spectroscopically) and reconfirmed with Keck + DEIMOS. This high redshift identification is enabled by the depth of the PEARS survey. Substantially higher redshifts are precluded for PEARS data by the declining sensitivity of the ACS grism at $\lambda > 0.95\mu\text{m}$. Spectra of Lyman breaks at yet higher redshifts will be possible using comparably deep observations with IR-sensitive grisms.

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1. Introduction

To properly understand the history of cosmic dawn, we must be able to reliably identify galaxies observed during the epoch of reionization. Such galaxies are the most likely sources of the radiation that ionized intergalactic hydrogen. They are the best places to look for signatures of primordial star formation, for even if the buildup of heavy elements is rapid,

the fraction of galaxies forming their first generations of stars should be higher if we observe them when the universe itself was young. The pace of their growth depends on incompletely understood physical processes — both the onset of star formation in low-metallicity conditions, and the potential disruption of later star formation by the ionizing radiation and/or supernovae produced by the first stellar generation. The best way to constrain the range of possible outcomes from these various processes is to take a direct, observational census of galaxies throughout the reionization era — from its end at $6 \lesssim z \lesssim 7$, back to the earliest galaxies we can identify.

Much progress has been made recently in this direction, thanks primarily to the dramatic increase in near-infrared imaging sensitivity and survey efficiency afforded by the Wide Field Camera 3 (WFC3) Infrared (IR) channel on the *Hubble Space Telescope* (*HST*). Imaging surveys with WFC3-IR have provided tens to hundreds of $z > 7$ galaxy candidates, identified by the Lyman- α absorption break in their broad band colors (e.g., Bouwens et al. 2010; Yan et al. 2010, 2011; Finkelstein et al. 2012). (We will refer to these as “Lyman break galaxies,” while noting that selection by a strong continuum break can identify either the 912Å break due to Lyman continuum absorption, or the 1216Å break due to Lyman- α absorption. Since the Lyman- α forest is optically thick for $z \gtrsim 5$, surveys for $z > 5$ galaxies use the Lyman- α absorption break, while those at $z \lesssim 3$ primarily identify the 912Å break.) These broad band *HST* searches have broken new ground primarily because the NIR sky is orders of magnitude darker in space. An alternative search method is narrow-band imaging, which can find Lyman- α emitting galaxies efficiently at selected redshifts ($z = 6.5, 6.9, 7.3, 7.7, 8.8$) where the line falls in dark windows in the night sky spectrum (e.g., Hu et al. 2002; Rhoads et al. 2004; Iye et al. 2006; Willis et al. 2008; Hu et al. 2010; Ouchi et al. 2010; Hibon et al. 2010; Tilvi et al. 2010; Kashikawa et al. 2011; Clément et al. 2012; Shibuya et al. 2011; Rhoads et al. 2012; Krug et al. 2012).

However, issues remain. Ground-based near-IR spectroscopy can only confirm these objects easily when they have strong Lyman- α lines in clean regions of the night sky spectrum. Thus, while dozens have been confirmed up to $z = 6.5$ (Hu et al. 2010; Ouchi et al. 2010; Kashikawa et al. 2011), only a handful are confirmed at higher redshifts (Iye et al. 2006; Rhoads et al. 2012; Shibuya et al. 2011; Pentericci et al. 2011; Ono et al. 2011; Schenker et al. 2011). The crucial Lyman- α line may be rare and/or weak at redshifts where the IGM was mostly neutral (and hence able to scatter Lyman- α photons). Meantime, sample contamination by foreground galaxies becomes an increasing worry at higher redshifts, where the volume available for such contaminants becomes large. Finally, the candidate lists from the highest redshift galaxy hunts can be disturbingly unstable, showing little overlap when different groups examine the same data, or even when the same group re-observes the same field (e.g. Yan et al. 2011; Oesch et al. 2012).

Slitless spectroscopy with the *Hubble Space Telescope* offers a solution to many of these issues. Space telescopes avoid the crippling effects of Earth’s atmosphere on the near-IR sky. *HST*’s spatial resolution is well matched to the sizes of high redshift galaxies. These slitless grisms thus provide unmatched sensitivity to continuum emission from faint, compact high redshift galaxies.

In this *Letter*, we present the highest redshift spectroscopic confirmation to date from *HST*’s slitless grisms: The galaxy PEARS-N-101687, which was identified based on its Lyman- α break at $z \approx 6.6$ in deep ACS G800L slitless spectra from the PEARS survey (“Probing Evolution and Reionization Spectroscopically”). While the sensitivity of the ACS grism declines at 9600, the WFC3 IR channel has a similar grism and can perform spectroscopic confirmations at even higher redshifts.

We organize the paper as follows. In section 2, we describe the PEARS survey observations and data analysis. In section 3, we present followup spectroscopy at higher spectral resolution from the Keck telescope. In section 4, we discuss the implications of our findings for the field of high-redshift galaxy hunting. We conclude in section 5. Most photometry discussed here is from the *HST*-GOODS survey (Giavalisco et al. 2004). We denote the GOODS filters F450W as “ B_{450} ,” F606W as “ V_{606} ,” F775W as “ i_{775} ,” and F850LP as “ z_{850} ,” and use the AB magnitude system. Throughout the paper, we adopt a Λ -CDM “concordance cosmology” with $\Omega_M = 0.27$, $\Omega_\Lambda = 0.73$, and $H_0 = 71 \text{ km s}^{-1} \text{ Mpc}^{-1}$ (see Spergel et al. 2007).

2. PEARS Grism Observations

PEARS is the most extensive systematic survey conducted with the G800L grism on the *HST*’s Advanced Camera for Surveys-Wide Field Camera (ACS-WFC). PEARS is an *HST* Treasury program led by S. Malhotra. It covers a total of nine fields, including one deep pointing in the *Hubble* Ultra Deep Field, and eight wide-field pointings (four each in the GOODS-North and GOODS-South regions). Each pointing was observed at three or four distinct roll angles, to mitigate the impact of overlap between spectra of nearby objects.

The *HST* slitless spectra were reduced using the aXe package (Kümmel et al. 2009), following closely the procedure used for the earlier GRISM ACS Program for Extragalactic Science (GRAPES) survey (Pirzkal et al. 2004). For each roll angle, the relative offsets of all exposures were determined using zero-order images and narrow emission lines. The data for each roll angle were ultimately combined into 2D spectroscopic stacks and extracted 1D spectra for each source and each observed position angle.

To identify and spectroscopically confirm the highest redshift Lyman break galaxies in

the survey, we followed a procedure based on Malhotra et al. (2005). We started with the GOODS v1.9 images and performed our own SExtractor photometry. We then applied a “liberal” i -dropout criterion to generate a list of candidate Lyman break galaxies. Since the GOODS data do not include observations redder than z band (and our candidate selection was done prior to the installation of WFC3), this ultimately amounts to using $i_{775} - z_{850} > 0.9$. For each of these objects, we calculated the net significance (“netsig”) parameter \mathcal{N} (Pirzkal et al. 2004) to determine which spectra might have sufficient information for a redshift measurement. (“Netsig” is defined by first sorting all pixels in a spectrum in descending order of signal-to-noise ratio; calculating the signal-to-noise ratio S_n obtained by combining flux from the brightest n pixels, for all n between 1 and the total number of pixels in the spectrum; and finally taking $\mathcal{N} = \max\{S_n\}$.)

After selecting candidates by $i - z$ color and ranking them by netsig, several PEARS team members (including SM, JER, NP, SC, & NG) examined the spectra by eye. We did this because a straight $i - z$ color cut can select extremely red objects (EROs) as well as Lyman break galaxies (LBGs), but the spectral signatures of the two are distinct. (EROs show a smooth rise towards the red, while LBGs at these redshifts show a step function at the redshifted wavelength of Lyman- α forest absorption.)

The highest redshift object identified in the PEARS-Wide fields through this process was the galaxy PEARS-N-101687, at equatorial coordinates 12:37:25.65 +62:17:43.5 (J2000). This object has magnitude $z_{850} = 26.16$, while it is undetected in the B_{450} , V_{606} , and i_{775} bands. Its grism redshift estimate is $z = 6.6 \pm 0.1$ based on the observed location of the Lyman α break. We show postage stamp images of the object from the GOODS data in Figure 1, and the 2D PEARS spectrum in Figure 2.

3. Keck Followup Observations

We selected several galaxies, including PEARS-N-101687, for followup observations during a Keck observing run of three half nights on UT 2007 April 13–15, using the DEep Imaging Multi-Object Spectrograph (DEIMOS; Faber et al. 2003). These observations were part of a program of a deep field followup led by Spinrad, Stern, and Dickinson, using Keck telescope time from the University of California system. PEARS-N-101687 was included on three slit masks during this run, and also on one additional mask observed in early 2008. All observations used the 600 line grism. The observations are summarized briefly in Table 1.

We show the extracted Keck + DEIMOS spectrum of PEARS-N-101687 in Figure 3. The galaxy shows a prominent Lyman- α line at redshift $z = 6.573$, consistent with the grism-

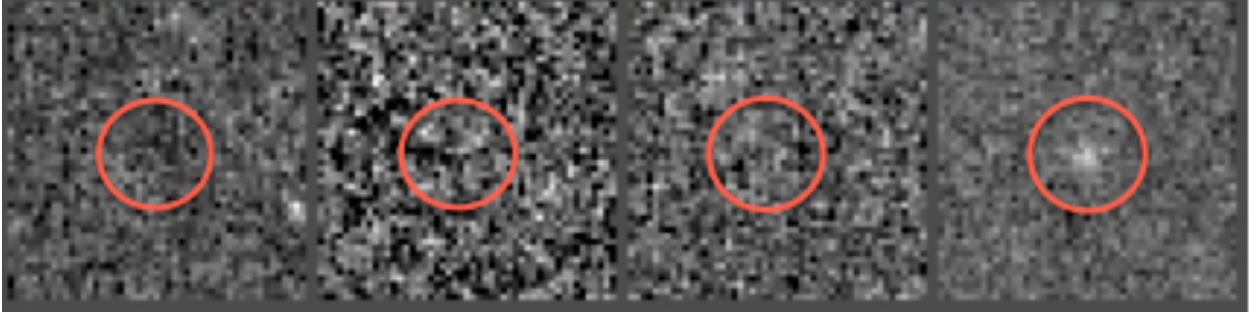


Fig. 1.— Direct imaging of PEARs-N-101687 from the GOODS survey (Giavalisco et al. 2004). The object is effectively undetected in the B_{450} , V_{606} , and i_{775} bands (left three panels), while it is clearly seen in the z_{850} band (right panel) with magnitude $z_{850} = 26.16$ (AB). Each panel is $1.5'' \times 1.5''$ in size.

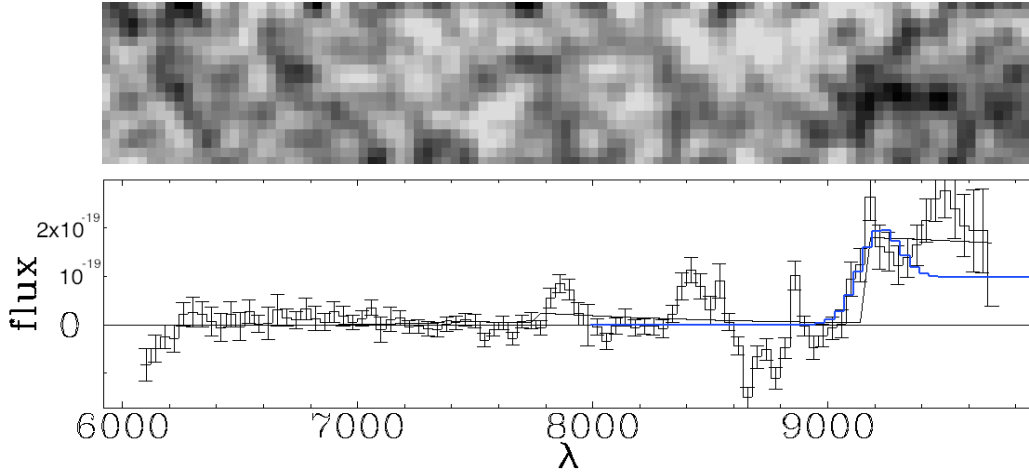


Fig. 2.— *Upper*: The 2D PEARs spectrum of PEARs-N-101687, displayed in inverse video (dark = more flux), with all three position angles coadded and with a 2 pixel ($0.1'' \times 80\text{\AA}$) FWHM smoothing applied. A short segment of continuum is evident, with a blue edge near 9230\AA due to the Lyman- α forest break, and a red cutoff imposed by the falloff of instrumental efficiency. *Lower*: A 1D extraction of the PEARs spectrum, in $\text{ergs cm}^{-2} \text{s}^{-1} \text{\AA}^{-1}$. The thin black curve is a best-fitting continuum + IGM absorption model with redshift $z = 6.6$ and a continuum level of $AB = 25.7$ mag on the red side of the Lyman- α break. The blue curve is a model based on the line flux from the Keck spectrum, and the continuum flux level and line spread function width expected based on *HST* imaging.

based redshift of $z = 6.6 \pm 0.1$. This line is plainly detected in each of the three masks that have > 1 hour of exposure time. Averaging the spectra from the three useful masks together yields a better detection of the asymmetric line, along with hints of the continuum on the red side of the line. No other emission features are convincingly present in the DEIMOS spectrum, which spans approximately $5000\text{\AA} < \lambda < 10000\text{\AA}$.

To estimate the spectroscopic line flux, we first calibrated the observed spectra using a DEIMOS sensitivity function derived for the same grism but on another night. To compensate for differences in throughput (e.g., due to differences in slit losses or atmospheric extinction between nights), we identified two objects in the 2008 mask with reasonably bright continuum flux ($i_{775} \sim 21\text{--}22$ mag). For each, we weighted the spectrum by the throughput of the *HST* ACS F775W filter, integrated, and compared the result to the broad band photometry from the GOODS project. This yielded a mean correction of about 30%, relative to our archival sensitivity function. The reference objects used are both galaxies at moderate redshift, and although they are not point sources, they are considerably smaller than the DEIMOS slit. We measured their fluxes in a virtual $1''$ slit in ACS images both before and after smoothing with a $1''$ Gaussian “seeing.” We thereby found that differential slit losses would require a 10% correction to the spectroscopic flux of a point source. Applying our calibration to the spectrum of PEARS-N-101687 (and assuming the source is pointlike in $1''$ seeing), we find a spectroscopic line flux of $f_{\text{Ly}\alpha} \approx (2.8 \pm 0.6) \times 10^{-17} \text{ erg cm}^{-2} \text{ s}^{-1}$, where the uncertainty is dominated by the flux calibration of the spectrum.

4. Discussion

4.1. Other observations of PEARS-N-101687

While no spectrum of PEARS-N-101687 appears to have been previously published, the object is listed as Lyman- α candidate in Hu et al. (2010) (table 3, second entry), based on a narrowband excess in a filter with 9210\AA central wavelength and 120\AA FWHM. The narrow-band magnitude published in that work, $AB = 24.36$, corresponds to a total flux of

Mask	Obs dates (UT)	$N_{\text{exp}} \times \text{Duration}$	Total time	Conditions	Comments
hdf07c	2007 Apr 14	$4 \times 1800\text{s}$	7200 s	clear, $> 1''$ seeing	Good
hdf07d	2007 Apr 14–15	$5 \times 1800\text{s}$	9000 s	clear, $> 1''$ seeing	Good
hdf07e	2007 Apr 16	2820s	2820s	clear, $1.4''$ seeing	Not useful
hdf08a	2008 Mar 06	$6 \times \sim 1800\text{s}$	10500 s		Good

Table 1: Log of DEIMOS observations of PEARS-N-101687.

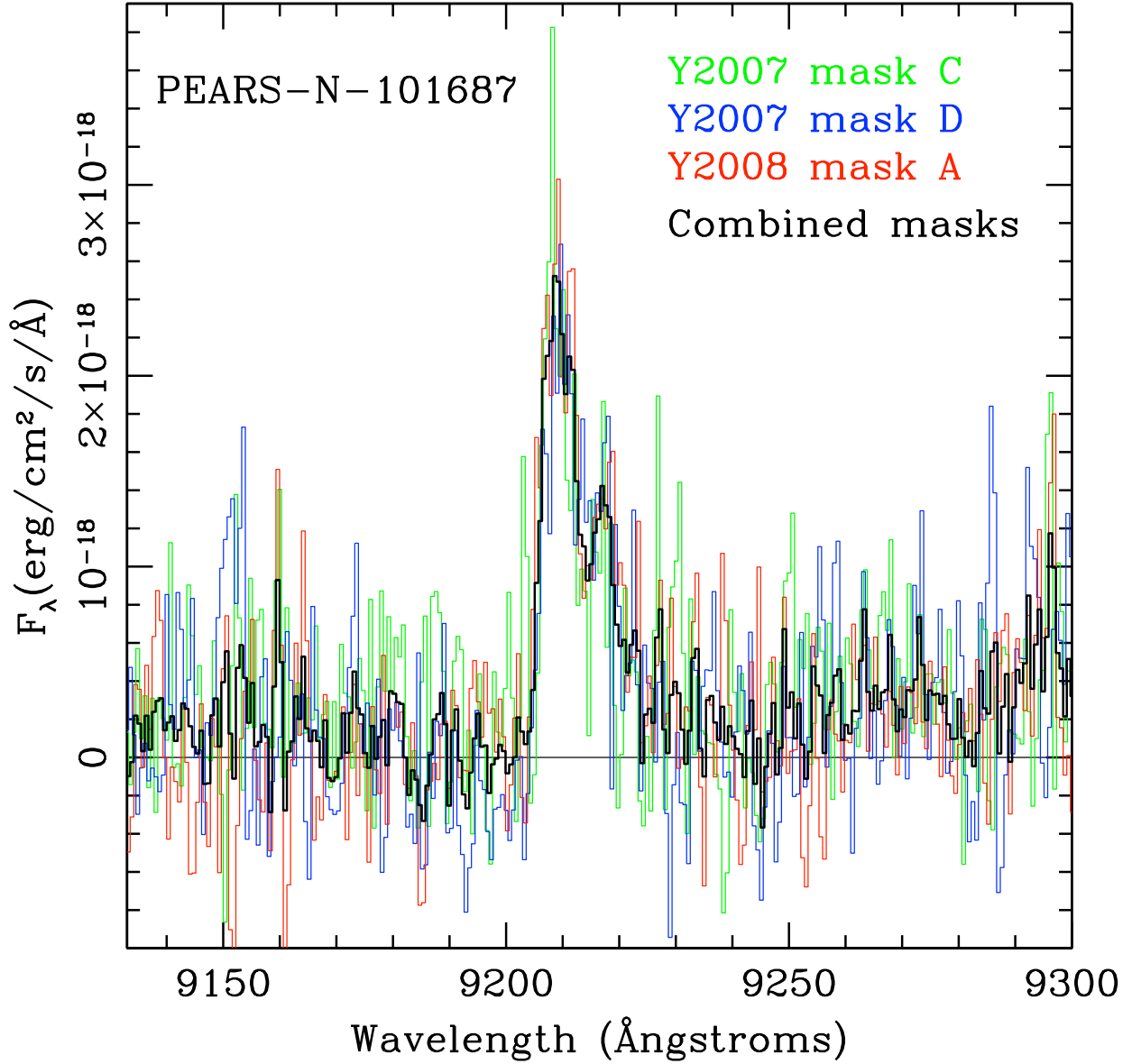


Fig. 3.— The 1D Keck + DEIMOS spectrum of PEARs-N-101687. The prominent line at 9210Å is identified as Lyman- α based on the strong break at 9210Å in the *HST* PEARs grism spectrum, the complete absence of flux blueward of the line, and the prominent asymmetry of the line. The spectra from each individual deep Keck+DEIMOS data set is shown with its own color, and the exposure-time weighted mean of these three masks is shown in black. The spectra are flux calibrated (with $\sim \pm 20\%$ accuracy) as described in the text. There is a suggestion of continuum flux on the red side of the line in the sum of all the Keck data. No other convincing features are seen in the full DEIMOS spectrum, which fully covers the range 5000–10000Å.

about $2.75 \times 10^{-17} \text{ erg cm}^{-2} \text{ s}^{-1}$ within the 120\AA filter.

The CANDELS survey (Grogin et al. 2011; Koekemoer et al. 2011) provides WFC3-IR photometry of this region. PEARS-N-101687 is detected at high confidence in the near infrared. CANDELS catalog magnitudes for the source are $z_{850} = 26.36 \pm 0.18$, $Y_{105} = 25.27 \pm 0.11$, $J_{125} = 25.11 \pm 0.09$, and $H_{160} = 25.04 \pm 0.08$. (Note the small [0.2 mag, 1σ] difference between our previous z_{850} and the CANDELS project photometry.) The Y_{105} magnitude is near the wavelength of the Lyman- α line but unaffected by Lyman- α forest absorption. Using the Y_{105} magnitude to estimate the continuum flux density just redward of the Lyman- α line, we expect $\approx 6 \times 10^{-18} \text{ erg cm}^{-2} \text{ s}^{-1}$ of continuum flux in the narrowband filter used by Hu et al. (2010). This leaves $\sim 2.1 \times 10^{-17} \text{ erg cm}^{-2} \text{ s}^{-1}$ as the line flux expected based on photometry. This is 25% below the flux we derive from the Keck Deimos observations, but consistent within the combined uncertainties of all the data sets involved.

4.2. Comparison of grism and slit spectra

The galaxy PEARS-N-101687 is reminiscent of the bright $z = 5.83$ dropout galaxy UDF2225 (Malhotra et al. 2005) = SiD2 (Dickinson et al. 2004) = SBM3 (Stanway et al. 2003, 2004), in that the ACS grism spectrum shows a clear continuum with a Lyman break, while the followup slit spectrum from Keck shows a prominent Lyman- α line. This contrast is a consequence of the differing capabilities of the two instruments. The high spatial resolution and low sky background of *HST* + ACS provide exquisite sensitivity to faint continuum emission. On the other hand, the higher spectral resolution of Keck+DEIMOS slit spectra provides a clearer look at the Lyman- α line. The PEARS ACS spectrum does have sufficient sensitivity to detect the observed Lyman- α line, but the line is blended with the Lyman break at the resolution of the grism. For the Lyman- α line to appear obvious at $z \gtrsim 5$, where the Lyman- α forest is optically thick, its observer-frame equivalent width must exceed the instrumental resolution, which is about $150\text{--}200\text{\AA}$ for PEARS-N-101687 (based on its half-light diameter in the GOODS images). The observed equivalent width is modestly larger than this threshold, but not so large that the line is expected to be prominent in the grism spectrum. The 2D HST spectrum also shows a hint of extended Lyman- α emission (see Rhoads et al. 2009; Bond et al. 2010; Finkelstein et al. 2011), which would lie outside the extraction region for PEARS 1D spectra but within the wider Keck slit, further increasing the relative prominence of the line in the Keck spectrum.

We have modeled the expected 1D grism spectrum by assuming a flat continuum at the level of the Y_{105} flux measurement, a Lyman- α line at the wavelength and flux observed by Keck, and a line spread function determined by the observed z_{850} angular size of the object

and the dispersion of the ACS G800L grism. The resulting model is shown as a blue curve in figure 2. Near the Lyman- α line, from 9000–9400Å, the agreement is quite good. The largest discrepancy is at yet redder wavelengths, where the grism spectrum continuum appears to be above the Y_{105} flux. This may be due to an edge effect always present in flux-calibrated slitless spectra of extended sources. The counts at a particular pixel include redder light from one edge of the source, and bluer light from the other edge, but all counts in the pixel are converted to flux density using a single system throughput, which is the one appropriate for light from the centroid of the source that is dispersed onto that pixel. Where the efficiency is changing rapidly with wavelength, as it does in the 9500Å region for this instrument, the net effect is an over-estimate of the flux. Bumps in the 1D ACS spectrum at 7850Å and 8400Å could be contamination by other fainter sources, transmissive gaps in the OH forest absorption, or simply regions of somewhat correlated noise in the extracted spectrum.

4.3. PEARS-N-101687 in context

The galaxy PEARS-N-101687 is, for its redshift, a moderately bright object. It has a 1500Å absolute magnitude $M_{1500} = -21.38$, based on an interpolation of the CANDELS Y_{105} and J_{125} fluxes. Compared to the published luminosity function for a sample of candidate $z \approx 6.6$ LBGs from Bouwens et al. (2011), PEARS-N-101687 is about $3\times$ (or 1.25 magnitudes) brighter than L^* (the characteristic galaxy luminosity for the best fit Schechter function).

The Lyman- α luminosity of PEARS-N-101687 is $L_{Ly\alpha} = 1.4 \times 10^{43} \text{ erg s}^{-1}$, based on its line flux from the DEIMOS spectrum and a “concordance” cosmology luminosity distance of $d_L = 65.5 \text{ Gpc}$. This is on the bright end of the distribution for $z = 6.5$ narrowband-selected Lyman- α galaxy samples, which yield Schechter function fits with characteristic luminosities of $L_* = 4.4 \times 10^{42}$, 5.8×10^{42} , and $1.0 \times 10^{43} \text{ erg s}^{-1}$, respectively, for Ouchi et al. (2010), Kashikawa et al. (2011), and Hu et al. (2010).

The spectroscopic line flux of PEARS-N-101687, combined with its Y_{105} flux density, yields an observer frame equivalent width of $EW = 290 \pm 80 \text{ Å}$, or in the rest frame, $EW_0 = 38 \pm 12 \text{ Å}$. This is below the average for narrowband-selected samples, as one might expect given that we identified it by its continuum trace in the PEARS spectrum. Lyman- α galaxies often have rest frame equivalent widths above 200Å (Malhotra & Rhoads 2002), and at $z = 6.5$, over 75% of the narrowband selected Lyman- α emitters have $EW_0 > 40 \text{ Å}$ (Kashikawa et al. 2011; Ouchi et al. 2010).

The Lyman- α line asymmetry in PEARS-N-101687 is prominent, even by the standards

of high redshift Lyman- α emitting galaxies. Using the asymmetry measures¹ from Rhoads et al. (2003), we find $a_\lambda = 3.59$ and $a_f = 3.47$. Comparing these to a sample of 58 Lyman- α emitters observed using the same spectrograph and grating (Dawson et al. 2007), PEARS-N-101687 has the largest value of a_f , and the fifth-largest value of a_λ . This may be partly due to Lyman- α forest absorption of flux on the blue side of the systemic velocity, since the Lyman- α forest optical depth at $z = 6.57$ will exceed that at the lower redshift ($z = 4.5$) sample of Dawson et al. (2007).

In addition to its asymmetry, the line profile shows a dip at 9214Å, separating a secondary peak (at 9217Å) from the primary one (at 9207Å). The dip is about 4Å wide; the flux density there drops by about 40% relative to a smoothed “envelope” of the line flux; and the feature is significant at about the 3σ level. Such a feature could be explained by neutral gas in front of the emitter; two emitting regions separated by about 330 km s⁻¹ in line-of-sight velocity; or some more complex interplay of Lyman- α emission and scattering in a moving medium.

4.4. The need for spectroscopy

Samples of well over 100 Lyman break selected galaxy candidates have now been published for photometric redshifts $z > 7$, based on the combination of deep photometry at optical and near-IR wavelengths using *HST* (e.g., Bouwens et al. 2010; Yan et al. 2010, 2011; Finkelstein et al. 2012). This represents substantial progress in understanding galaxy evolution around the end of the reionization era. *However, essentially all these objects remain candidates at the moment, unconfirmed by spectroscopy.* Deep grism spectra from *HST* can play a unique role in fixing this.

The overlap between samples published by different groups can be distressingly small, even when those groups use exactly the same data sets. For example, consider recent publications on the bright end of the $z \approx 8$ LBG luminosity function by Yan et al. (2011) and Oesch et al. (2012), both using the first epoch of the CANDELS deep observations of the GOODS-S region. The two groups publish 8 and 9 candidate galaxies, respectively. However, only two objects are identified by both papers. Similar levels of inconsistency have been frequent in earlier studies. Indeed, a recent paper by Ellis et al. (2013) reports that *no* previously published galaxy candidate at $8.5 < z < 10$ in the Hubble Ultra Deep Field

¹These are $a_\lambda = (\lambda_{10,r} - \lambda_p)/(\lambda_p - \lambda_{10,b})$ and $a_f = \int_{\lambda_p}^{\lambda_{10,r}} f_\lambda d\lambda / \int_{\lambda_{10,b}}^{\lambda_p} f_\lambda d\lambda$. Here λ_p is the wavelength where the line peaks, and $\lambda_{10,b}$ and $\lambda_{10,r}$ are the wavelengths where the flux falls to 10% of peak on the blue and red sides of the line.

remains a viable high-redshift object after the addition of deeper imaging in the WFC3-IR F105W (“ Y_{105} ”) and F140W (“ JH_{140} ”) filters, while reporting a set of seven new candidates in that redshift range.

Several factors contribute to unreliable candidate lists. First, candidates are generally sought down to the limit of the survey depth, and objects near the faint limit of the data will then inevitably outnumber brighter, better measured sources. This means that photometric noise can push a galaxy across the selection line — in *or* out of a candidate sample, either in brightness or in color. Apparently minor differences in the choice of photometry method (apertures of various radii vs. SExtractor “magauto”; different methods of sky background estimation; etc) can thus change samples appreciably.

Additionally, different authors may choose somewhat different criteria in selecting their candidates. Some use photometric redshifts (Finkelstein et al. 2012), while most others use straight color and magnitude cuts (e.g., Bouwens et al. 2010; Oesch et al. 2012; Yan et al. 2011), but the adopted cuts are not always the same.

As the search redshift increases, the volume of foreground space and the variety of possible foreground contaminants increases too. For galaxies at $z > 7$, plausible foreground contaminants include Galactic brown dwarfs, and both early-type galaxies and ultra-strong emission line sources at intermediate redshifts (e.g. the candidate lensed $z = 11$ galaxy A2667-J1, whose spectrum revealed it to be an [OIII] $\lambda\lambda 4959, 5007$ emitter at $z = 2.082$; Hayes et al. 2012). While published $z > 7$ candidates usually have spectral energy distributions that are less well fit by any foreground model than by a Lyman break galaxy at $z > 7$, a majority allow viable $z \ll 7$ solutions (Pirzkal et al. 2013, in prep), and the number of foreground objects in the survey area may exceed the number of $z > 7$ galaxies substantially. *Incorporating this as prior information, the fraction of $z > 7$ candidates that are actual $z > 7$ galaxies may be modest.*

Spectroscopic followup of $z > 7$ candidates can resolve these uncertainties. Ground-based spectra can confirm true $z > 7$ galaxies, but generally only when they have strong Lyman- α emission that is neither blocked by atmospheric H₂O absorption nor blended with strong OH airglow lines at the resolution of the spectrograph. Likewise, ground-based spectra may definitively rule out foreground objects whose Lyman-break colors are due to strong emission lines in one or two filters.

However, sources without strong emission lines require continuum spectroscopy for definitive confirmation. For redshifts $z < 6$, a minority of LBGs have strong Lyman- α emission (Steidel et al. 2000; Stark et al. 2010). As we push to higher redshifts, within the epoch of reionization, Lyman- α will be obscured by resonant scattering in an increasingly

neutral intergalactic medium (IGM) (Miralda-Escude & Rees 1998; Haiman & Spaans 1999; Rhoads & Malhotra 2001). This effect offers valuable tests of reionization (e.g., Malhotra & Rhoads 2004; Stern et al. 2005; Pentericci et al. 2011; Ono et al. 2012; Schenker et al. 2012). It also means that *continuum break spectroscopy is crucial* for spectroscopic confirmations in the epoch of neutral gas.

Such spectroscopy is impractical from the ground with current instruments. From space, the absence of OH emission lines and water absorption makes the job much easier. To obtain a Lyman break confirmation at the knee of the luminosity function (L^*), we need to detect the continuum with good statistical significance, over a wavelength range of several hundred Ångstrom on the red side of the Lyman break. Deep integrations with *HST* grisms can accomplish this, spectroscopically confirming redshifts for objects as faint as 27th magnitude (this work; Malhotra et al. 2005; Rhoads et al. 2009).

5. Summary/Conclusions

We present here PEARS-N-101687, the highest redshift galaxy identified in the wide-field component of the PEARS slitless spectroscopic survey, which achieved a depth of 20 orbits of *HST* ACS G800L slitless spectroscopy over 80 square arcminutes. This galaxy, with a grism redshift of $z = 6.6 \pm 0.1$, remains the highest redshift object yet confirmed with *HST* slitless spectroscopy. While fainter objects could be spectroscopically confirmed at the sensitivity limit of our ACS slitless spectroscopy in the Hubble Ultra Deep Field — where the aggregate depth approaches 90 orbits — the wavelength limit of the ACS Wide Field Camera CCD detectors effectively limits the survey redshifts to $z \lesssim 6.7$ (e.g., Malhotra et al. 2005).

Keck telescope + DEIMOS followup of this object confirms and refines its grism redshift: $z = 6.57$. The object has a rest-frame UV continuum magnitude $M_{1500} = -21.38$, a Lyman- α line luminosity of $1.4 \times 10^{43} \text{ erg s}^{-1}$, and a rest-frame equivalent width of $38 \pm 12 \text{ Å}$. This makes it a relatively luminous Lyman- α emitting Lyman break galaxy. We emphasize that the discovery and redshift from *HST* PEARS are based primarily on the continuum and Lyman- α forest break, and not on the Lyman- α emission line, which is not prominent in the *HST* spectrum.

The discovery of this object demonstrates the value of deep continuum observations with *HST* slitless grisms for spectroscopic confirmation of galaxies in the epoch of reionization. Comparably sensitive observations with the *HST* WFC3-IR channel grisms have the potential to provide Lyman break spectroscopic confirmations of $z > 7$ galaxies — something

that still eludes our other observational capabilities, and that now presents a large obstacle in advancing our understanding of galaxy evolution in the era of cosmic dawn.

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REFERENCES

- Bond, N. A., Feldmeier, J. J., Matković, A., Gronwall, C., Ciardullo, R., & Gawiser, E. 2010, *ApJ*, 716, L200
- Bouwens, R. J., et al. 2010, *ApJ*, 709, L133
- . 2011, *ApJ*, 737, 90
- Clément, B., et al. 2012, *A&A*, 538, A66
- Dawson, S., Rhoads, J. E., Malhotra, S., Stern, D., Wang, J., Dey, A., Spinrad, H., & Jannuzi, B. T. 2007, *ApJ*, 671, 1227
- Dickinson, M., et al. 2004, *ApJ*, 600, L99
- Ellis, R. S., et al. 2013, *ApJ*, 763, L7

- Faber, S. M., et al. 2003, in Society of Photo-Optical Instrumentation Engineers (SPIE) Conference Series, Vol. 4841, Society of Photo-Optical Instrumentation Engineers (SPIE) Conference Series, ed. M. Iye & A. F. M. Moorwood, 1657–1669
- Finkelstein, S. L., et al. 2011, *ApJ*, 735, 5
- . 2012, *ApJ*, 756, 164
- Giavalisco, M., et al. 2004, *ApJ*, 600, L93
- Grogin, N. A., et al. 2011, *ApJS*, 197, 35
- Haiman, Z., & Spaans, M. 1999, *ApJ*, 518, 138
- Hayes, M., Laporte, N., Pelló, R., Schaerer, D., & Le Borgne, J.-F. 2012, *MNRAS*, 425, L19
- Hibon, P., et al. 2010, *A&A*, 515, A97
- Hu, E. M., Cowie, L. L., Barger, A. J., Capak, P., Kakazu, Y., & Trouille, L. 2010, *ApJ*, 725, 394
- Hu, E. M., Cowie, L. L., McMahon, R. G., Cpaak, P., Iwamuro, F., Kneib, J.-P., Maihara, T., & Motohara, K. 2002, *ApJ*, 568, L75
- Iye, M., et al. 2006, *Nature*, 443, 186
- Kashikawa, N., et al. 2011, *ApJ*, 734, 119
- Koekemoer, A. M., et al. 2011, *ApJS*, 197, 36
- Krug, H. B., et al. 2012, *ApJ*, 745, 122
- Kümmel, M., Walsh, J. R., Pirzkal, N., Kuntschner, H., & Pasquali, A. 2009, *PASP*, 121, 59
- Malhotra, S., & Rhoads, J. E. 2002, *ApJ*, 565, L71
- . 2004, *ApJ*, 617, L5
- Malhotra, S., et al. 2005, *ApJ*, 626, 666
- Miralda-Escude, J., & Rees, M. J. 1998, *ApJ*, 497, 21
- Oesch, P. A., et al. 2012, arXiv, astro-ph.CO
- Ono, Y., et al. 2011, arXiv, astro-ph.CO

- Ono, Y., et al. 2012, ApJ, 744, 83
- Ouchi, M., et al. 2010, ApJ, 723, 869
- Pentericci, L., et al. 2011, ApJ, 743, 132
- Pirzkal, N., et al. 2013, in prep
- Pirzkal, N., et al. 2004, ApJS, 154, 501
- Rhoads, J. E., Hibon, P., Malhotra, S., Cooper, M., & Weiner, B. 2012, ApJ, 752, L28
- Rhoads, J. E., & Malhotra, S. 2001, ApJ, 563, L5
- Rhoads, J. E., et al. 2003, AJ, 125, 1006
- . 2004, ApJ, 611, 59
- . 2009, ApJ, 697, 942
- Schenker, M. A., Stark, D. P., Ellis, R. S., Robertson, B. E., Dunlop, J. S., McLure, R. J., Kneib, J. P., & Richard, J. 2011, arXiv, astro-ph.CO
- Schenker, M. A., Stark, D. P., Ellis, R. S., Robertson, B. E., Dunlop, J. S., McLure, R. J., Kneib, J.-P., & Richard, J. 2012, ApJ, 744, 179
- Shibuya, T., Kashikawa, N., Ota, K., Iye, M., Ouchi, M., Furusawa, H., Shimasaku, K., & Hattori, T. 2011, eprint arXiv, 1112.3997, 11 pages, 11 figures, submitted to ApJ
- Spergel, D. N., et al. 2007, ApJS, 170, 377
- Stanway, E. R., Bunker, A. J., & McMahon, R. G. 2003, MNRAS, 342, 439
- Stanway, E. R., Bunker, A. J., McMahon, R. G., Ellis, R. S., Treu, T., & McCarthy, P. J. 2004, ApJ, 607, 704
- Stark, D. P., Ellis, R. S., Chiu, K., Ouchi, M., & Bunker, A. 2010, MNRAS, 408, 1628
- Steidel, C. C., Adelberger, K. L., Shapley, A. E., Pettini, M., Dickinson, M., & Giavalisco, M. 2000, ApJ, 532, 170
- Stern, D., Yost, S. A., Eckart, M. E., Harrison, F. A., Helfand, D. J., Djorgovski, S. G., Malhotra, S., & Rhoads, J. E. 2005, ApJ, 619, 12
- Tilvi, V., et al. 2010, arXiv, astro-ph.CO, submitted to ApJ, 3 figures

Willis, J. P., Courbin, F., Kneib, J.-P., & Minniti, D. 2008, MNRAS, 384, 1039

Yan, H., et al. 2011, eprint arXiv, 1112, 6406, submitted to ApJL

Yan, H.-J., Windhorst, R. A., Hathi, N. P., Cohen, S. H., Ryan, R. E., O’Connell, R. W., & McCarthy, P. J. 2010, Res. Astron. Astrophys., 10, 867